

Evaluation of Coefficients

Step 1. Calculate  $N_r^m$ 's from

21.7.8

$$N_{r+2}^m = \gamma_r^m - \lambda_{mn} - \frac{\beta_r^m}{N_r^m} \quad (r \geq 2)$$

$$N_2^m = \gamma_0^m - \lambda_{mn}; N_3^m = \gamma_1^m - \lambda_{mn}$$

$$\gamma_r^m = (m+r)(m+r+1)$$

$$+\frac{1}{2} c^2 \left[ 1 - \frac{4m^2 - 1}{(2m+2r-1)(2m+2r+3)} \right] \quad (r \geq 0)$$

Step 2. Calculate ratios  $\frac{d_0}{d_{2r}}$  and  $\frac{d_1}{d_{2p+1}}$  from

21.7.9 
$$\frac{d_0}{d_{2r}} = \left(\frac{d_0}{d_2}\right) \left(\frac{d_2}{d_4}\right) \dots \left(\frac{d_{2r-2}}{d_{2r}}\right)$$

21.7.10 
$$\frac{d_1}{d_{2p+1}} = \left(\frac{d_1}{d_3}\right) \left(\frac{d_3}{d_5}\right) \dots \left(\frac{d_{2p-1}}{d_{2p+1}}\right)$$

and the formula for  $N_r^m$  in 21.7.7.

The coefficients  $d_r^m$  are determined to within the arbitrary factor  $d_0$  for  $r$  even and  $d_1$  for  $r$  odd. The choice of these factors depends on the normalization scheme adopted.

Normalization of Angular Functions

Meixner-Schärfke Scheme

21.7.11 
$$\int_{-1}^1 [S_{mn}(c, \eta)]^2 d\eta = \frac{2}{2n+1} \frac{(n+m)!}{(n-m)!}$$

Stratton-Morse-Chu-Little-Corbató Scheme

21.7.12 
$$\sum_{r=0,1} \frac{(r+2m)!}{r!} d_r = \frac{(n+m)!}{(n-m)!}$$

(This normalization has the effect that  $S_{mn}(c, \eta) \rightarrow P_n^m(\eta)$  as  $\eta \rightarrow 1$ .)

Flammer Scheme [21.4]

21.7.13

$$S_{mn}(c, 0) = P_n^m(0) = \frac{(-1)^{\frac{n-m}{2}} (n+m)!}{2^n \left(\frac{n-m}{2}\right)! \left(\frac{n+m}{2}\right)!}$$

( $n-m$ ) even

21.7.14

$$S'_{mn}(c, 0) = P_n^{m'}(0) = \frac{(-1)^{\frac{n-m-1}{2}} (n+m+1)!}{2^n \left(\frac{n-m-1}{2}\right)! \left(\frac{n+m+1}{2}\right)!}$$

( $n-m$ ) odd

The above lead to the following conditions for

$d_r^m$

21.7.15

$$\sum_{r=0}^{\infty} \frac{(-1)^{r/2} (r+2m)!}{2^r \left(\frac{r}{2}\right)! \left(\frac{r+2m}{2}\right)!} d_r^m = \frac{(-1)^{\frac{n-m}{2}} (n+m)!}{2^{n-m} \left(\frac{n-m}{2}\right)! \left(\frac{n+m}{2}\right)!}$$

( $n-m$ ) even

21.7.16

$$\sum_{r=1}^{\infty} \frac{(-1)^{\frac{r-1}{2}} (r+2m+1)!}{2^r \left(\frac{r-1}{2}\right)! \left(\frac{r+2m+1}{2}\right)!} d_r^m = \frac{(-1)^{\frac{n-m-1}{2}} (n+m+1)!}{2^{n-m} \left(\frac{n-m-1}{2}\right)! \left(\frac{n+m+1}{2}\right)!}$$

( $n-m$ ) odd

(The normalization scheme 21.7.13 and 21.7.14 is also used in [21.10].)

Asymptotic Expansions for  $S_{mn}(c, \eta)$

21.7.17

$$S_{mn}(c, \eta) = (1-\eta^2)^{\frac{1}{2}} U_{mn}(c, \eta) \quad (c \rightarrow \infty)$$

$$U_{mn}(x) = \sum_{r=-\infty}^{\infty} h_r^l D_{l+r}(x) \quad l = n-m$$

where the  $D_r(x)$ 's are the parabolic cylinder functions (see chapter 19).

$$D_r(x) = (-1)^r e^{x^2/4} \frac{d^r}{dx^2} e^{-x^2/2} = 2^{-r/2} e^{-x^2/4} H_r\left(\frac{x}{\sqrt{2}}\right)$$

and the  $H_r(x)$  are the Hermite polynomials (see chapter 22). (For tables of  $h_{\pm r}^l/h_0^l$  see [21.4].)

Expansion of  $S_{mn}(c, \eta)$  in Powers of  $\eta$

21.7.18

$$S_{mn}(c, \eta) = (1-\eta^2)^{m/2} \sum_{r=0,1}^{\infty} p_r^{mn}(c) \eta^r$$

$$(r+1)(r+2)p_{r+2}^{mn}(c) - [r(r+2m+1) + m(m+1) - \lambda_{mn}(c)]p_r^{mn}(c) - c^2 p_{r-2}^{mn}(c) = 0$$

(The derivation of the transcendental equation for  $\lambda_{mn}$  is similar to the derivation of 21.7.4 from 21.7.3.)

Expansion of  $S_{mn}(c, \eta)$  in Powers of  $(1-\eta^2)$

21.7.19

$$S_{mn}(c, \eta) = (1-\eta^2)^{m/2} \sum_{k=0}^{\infty} c_{2k}^{mn} (1-\eta^2)^k \quad (n-m) \text{ even}$$